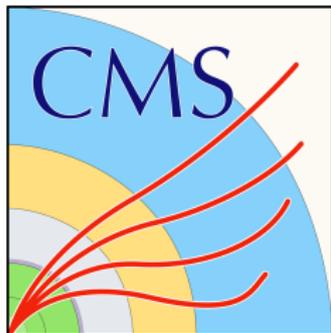


Dark Matter Analysis at Florida Tech

Merrick Lavinsky, Erick Yanes, Scott Demarest, Bandar Alsufyani

Florida Institute of Technology

3/4/26



Introduction to Dark Matter

- We don't know what most of the universe is made of!
- Dark matter is some invisible entity that warps spacetime
- Evident in galactic rotation curves, galaxy collisions/interactions
- Muon g-2 (Must be more forces), LHC Exotic matter (quark matter), Hubble constant tension

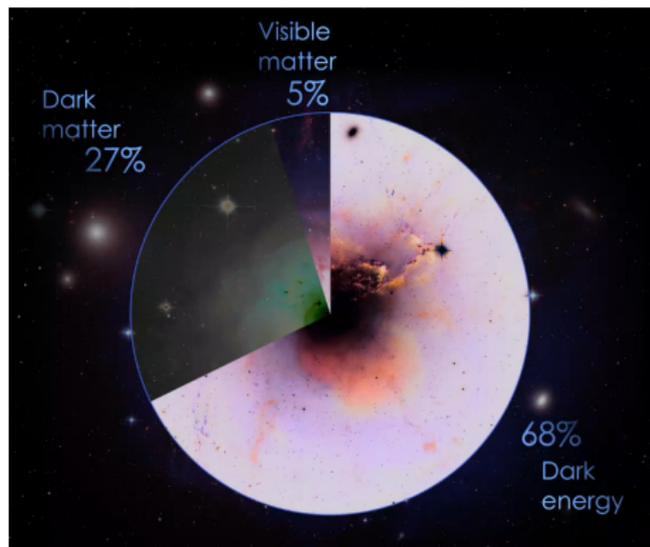


Figure: Current* measured composition of the universe via the Planck telescope [1]

Introduction to Dark Matter

But what *is* DM?

Introduction to Dark Matter

But what *is* DM?

Two groups: **new particle?**

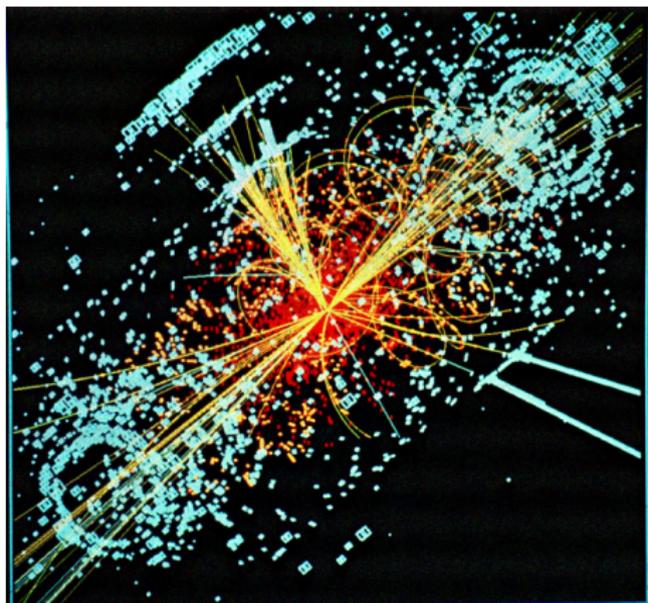


Figure: [2]

Introduction to Dark Matter

But what *is* DM?

Two groups: **new particle?**

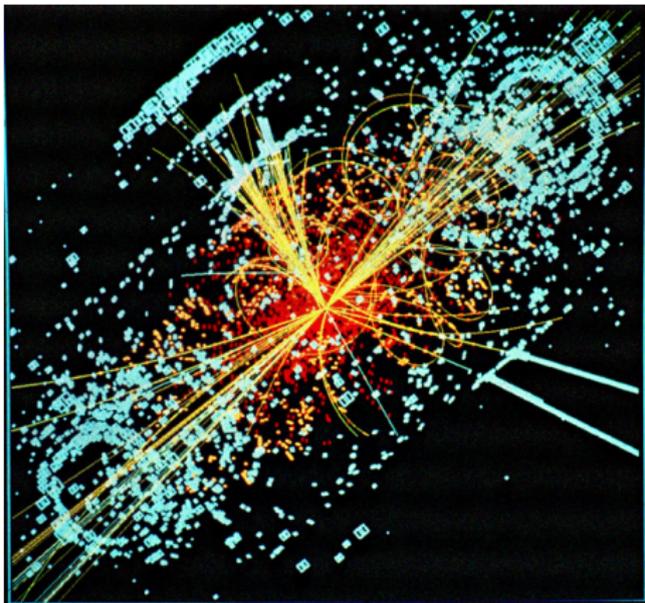


Figure: [2]

Modified Newtonian Dynamics (MOND)

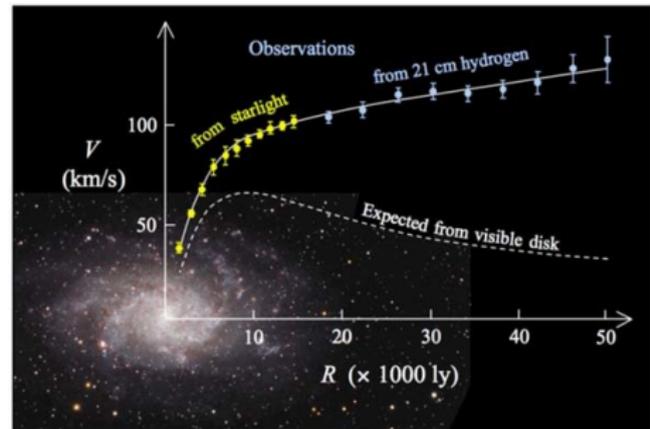


Figure: [3]

The CMS Detector and Particle Physics Lingo

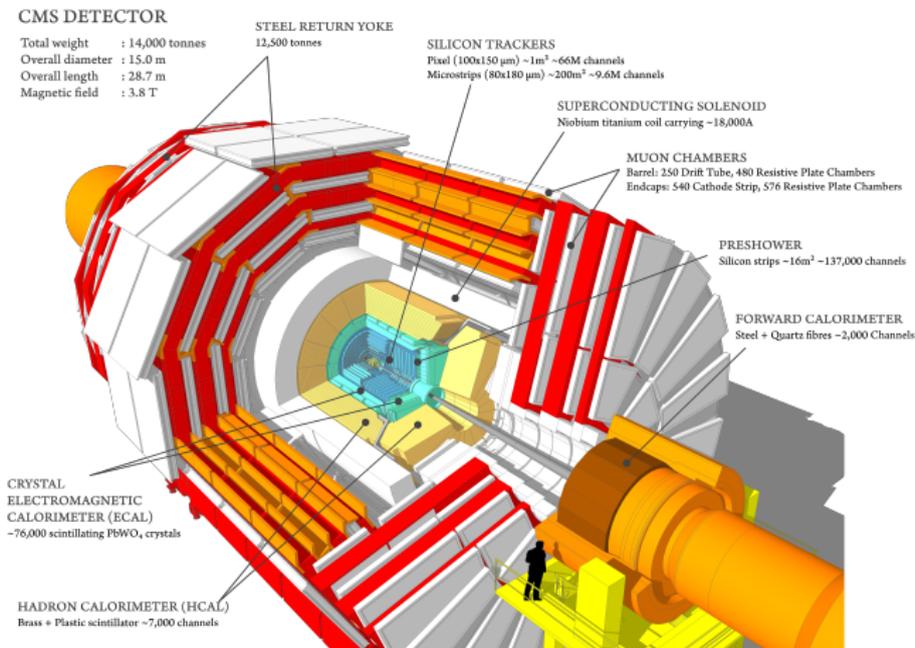


Figure: CMS Experiment at LHC. Collisions occur in the center and spread out through hundreds of subdetectors [4]

The CMS Detector and Particle Physics Lingo

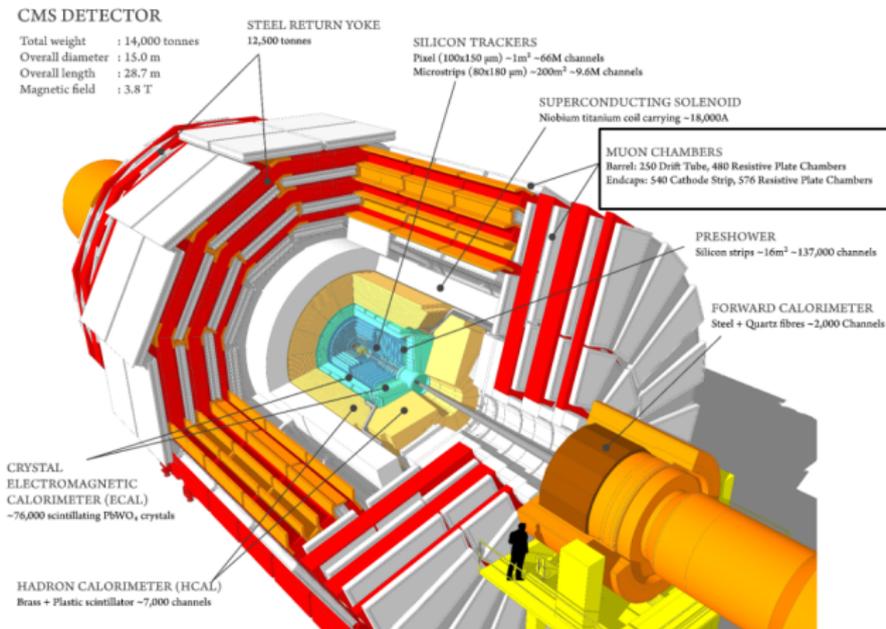


Figure: Muon Subsystem in outer barrel and endcap measures radiation coordinates and momentum of muons [4]

The CMS Detector and Particle Physics Lingo

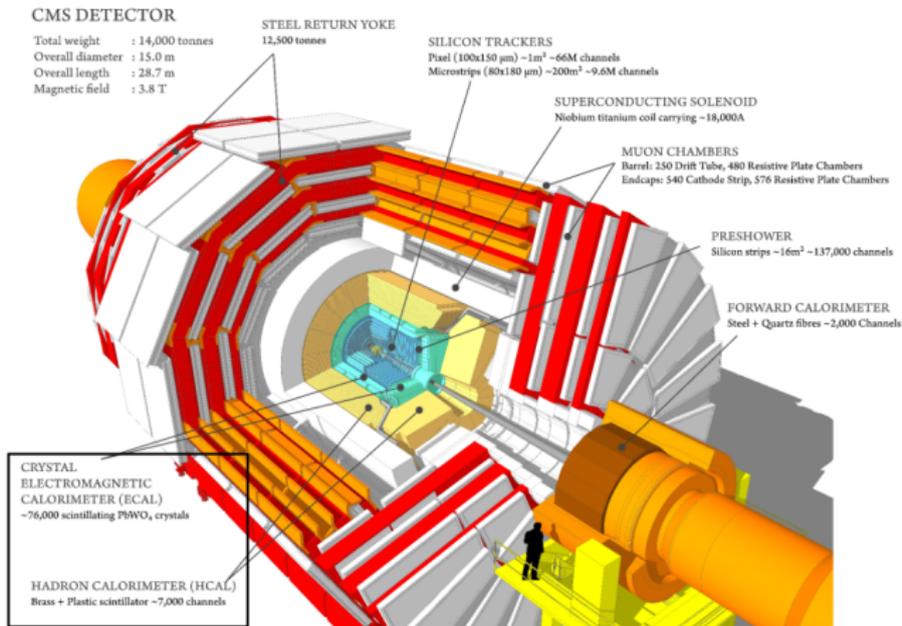


Figure: Electron Calorimeter (ECAL) and Hadron Calorimeter (HCAL) Subsystem in inner barrel to measure energy of electrons, photons, and light hadrons (mesons/baryons) [4]

The CMS Detector and Particle Physics Lingo

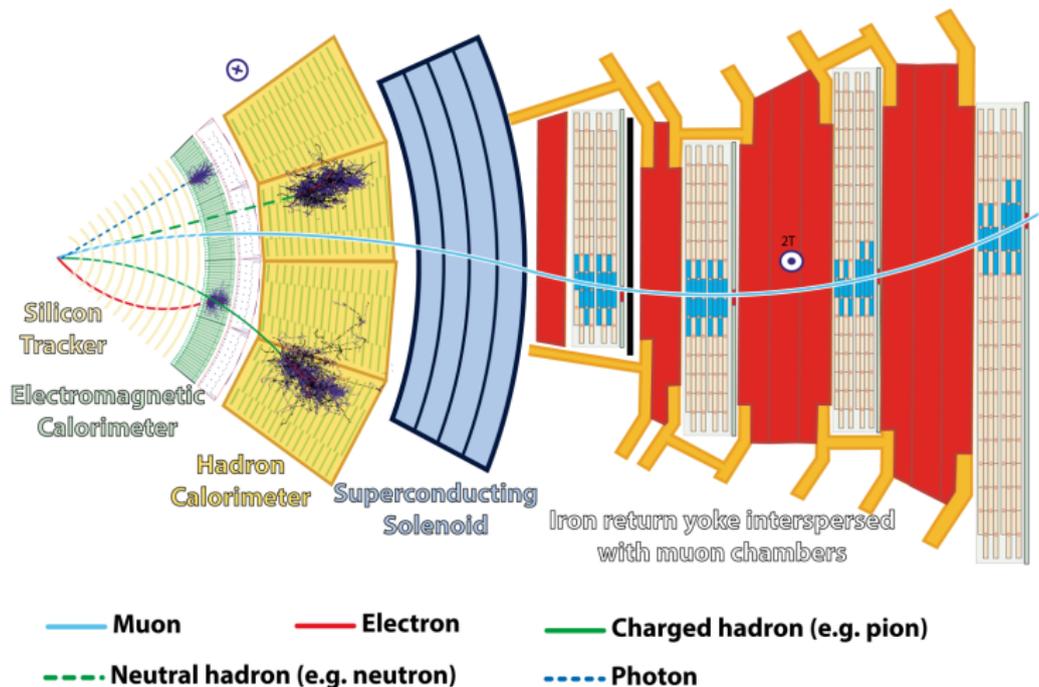


Figure: The various subsystems are designed to react with specific particles (hadrons/leptons/bosons) [5]

The CMS Detector and Particle Physics Lingo

- Cylindrical detectors \rightarrow cylindrical coordinates (ϕ, η, z)
- Relativistic QM \rightarrow Lorentz invariant coordinates = easier math

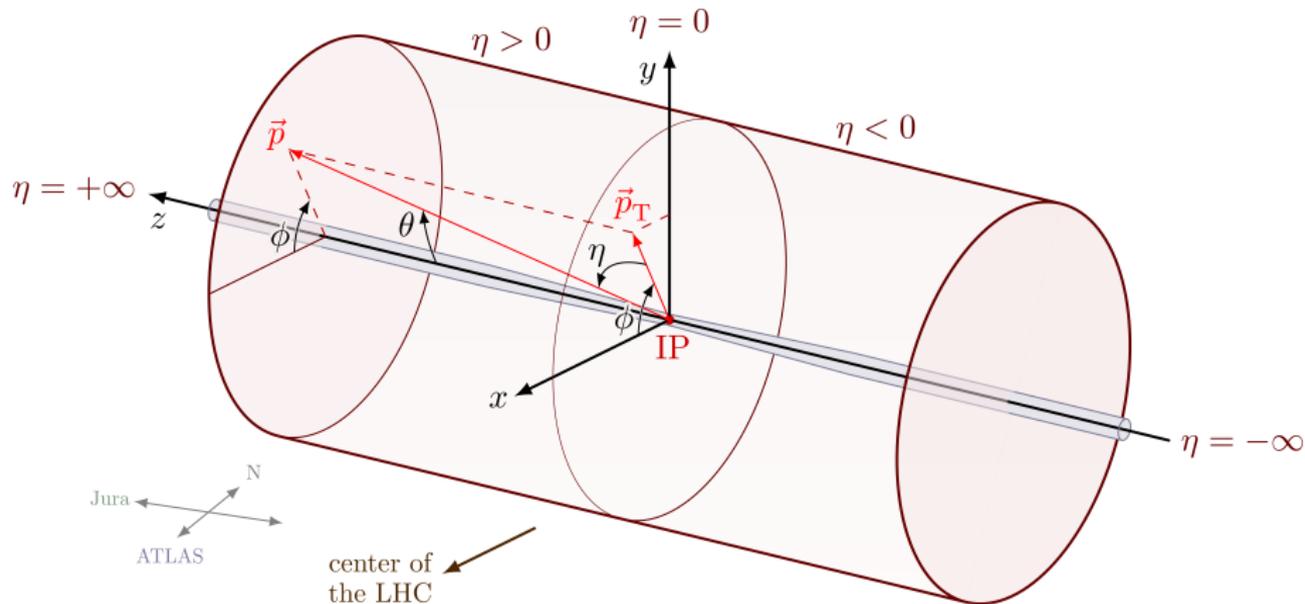


Figure: CMS detector coordinate system [6]

The CMS Detector and Particle Physics Lingo

- Relevant boundaries for CMS detector subsystems
- Nothing detected past $|\eta| \gtrsim 3$
- $1.44 \leq \eta \leq 1.57$ gap in calorimeters (Poor reconstruction)
- Muon reconstruction requires $|\eta| < 2.4$

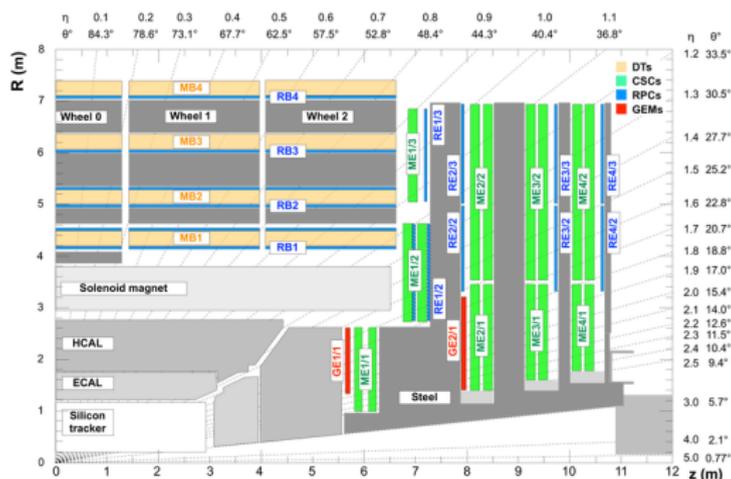
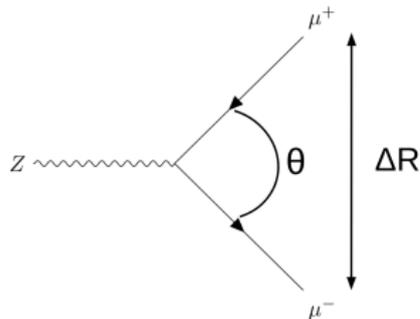


Figure: CMS detector quadrant with coordinate system interlayed [7]

The CMS Detector and Particle Physics Lingo

Particle decay

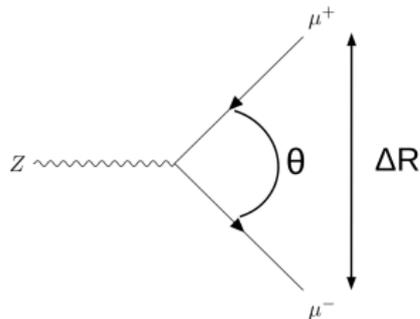
- $M_{\text{inv}} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2}$
- $\Delta R = \sqrt{(\eta_1 - \eta_2)^2 + (\phi_1 - \phi_2)^2}$
 - Angular distance
- $\theta \propto (M_P - m), P_P$
- Large mass to two small masses
→ highly collimated



The CMS Detector and Particle Physics Lingo

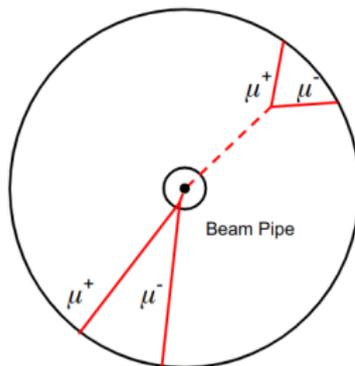
Particle decay

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Long Lived particles

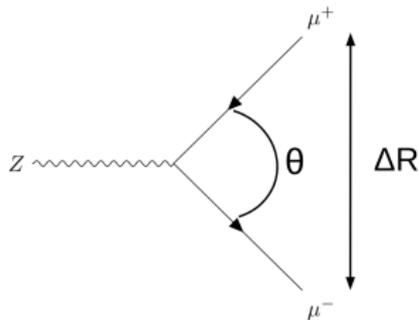
- $c\tau$ = decay length
- Central vertices = most common
- Dark matter = no interaction!
 - Large $c\tau$



The CMS Detector and Particle Physics Lingo

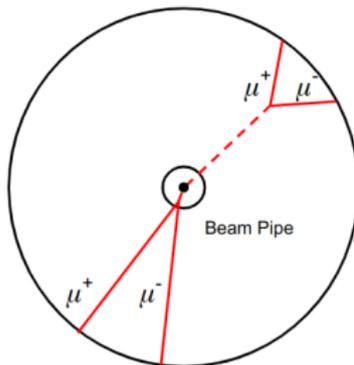
Particle decay

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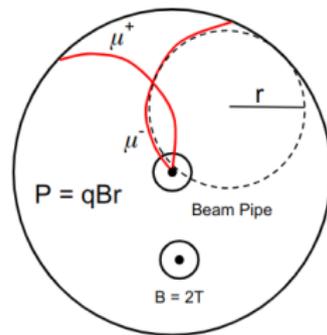
Long Lived particles

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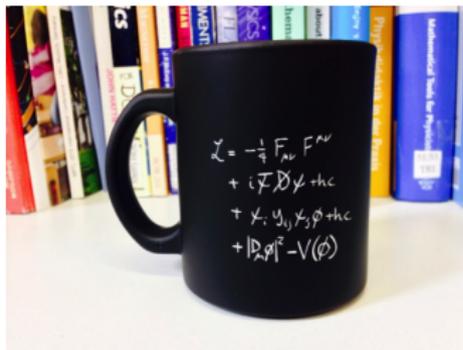


Momentum Calculation

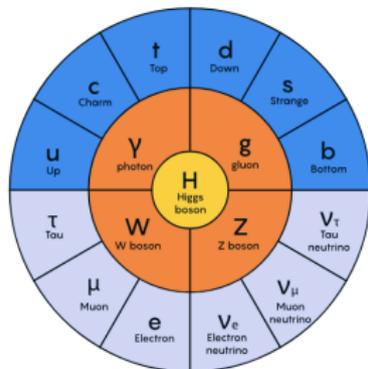
- Large B field (Solenoid)
- charged particles make S shape
- Momentum measured from curvature
- $P = qBr$



The CMS Detector and Particle Physics Lingo



The Standard Model



● QUARKS ● LEPTONS ● GAUGE BOSONS ● HIGGS BOSON

- Standard Model (SM) Lagrangian
 - Describes the fundamental forces and their interactions
 - $-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \rightarrow$ Electromagnetism (F), Weak Force (B), Strong Force (G)
 - $V(\phi) \rightarrow$ Higgs potential
- Most precise theory
- Bosons \rightarrow Force carriers
- Fermions \rightarrow Quarks, Leptons

[8]

HAHM basics

- Hidden Abelian Higgs Model (HAHM)
- What if *some* particles can produce DM
- Simplest extension to SM to include particles that could behave like DM
- Introduces dark Higgs and dark Z boson
- Free parameters: state mixing and masses
 - ϵ , M_{zd} - κ , M_{hd}

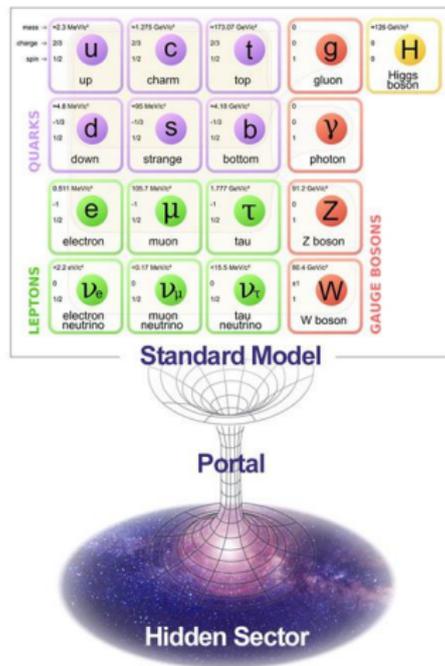


Figure: Standard model with a portal to a dark SM [9]

$$\mathcal{L}_{SM} = \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{1}{4} G_{\alpha}^{\mu\nu} G_{\mu\nu}^{\alpha} + \dots + V_0$$

- Start with the Standard Model
- Transform Hyperweak field: $B \rightarrow B^{\mu\nu} + \frac{\epsilon}{\cos(\theta_W)} Z^{\mu\nu}$
- Similar in H field
- New kinetic terms controlled by ϵ and κ

$$\mathcal{L}_{HAHM} = \frac{1}{4} Z^{\mu\nu} Z_{\mu\nu} + \frac{\epsilon}{2 \cos(\theta_W)} B^{\mu\nu} Z_{\mu\nu} + \frac{1}{2} m_{zd}^2 Z^{\mu\nu} Z_{\mu\nu}$$

$$V_0(H, S) = -\mu^2 |H|^2 + \lambda |H|^4 - \mu_S^2 |S|^2 + \lambda_S |S|^4 + \kappa |S|^2 |H|^2$$

So basically,

$$\begin{pmatrix} Z_{D,0} \\ B \end{pmatrix} = \begin{pmatrix} \sqrt{1 - \frac{\epsilon^2}{\cos^2 \theta}} & 0 \\ -\frac{\epsilon}{\cos \theta} & 1 \end{pmatrix} \begin{pmatrix} \hat{Z}_D \\ \hat{B} \end{pmatrix}$$

So basically,

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$$\Gamma(h \rightarrow Z_D Z_D) = \kappa'^2 \frac{1}{32\pi} \frac{v^2}{m_h} \sqrt{1 - \frac{4m_{Z_D}^2}{m_h^2}} \frac{(m_h^2 + 2m_{Z_D}^2)^2 - 8(m_h^2 - m_{Z_D}^2)m_{Z_D}^2}{m_h^4}$$

So basically,

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$$m_{Z,Z_D}^2 = \frac{1}{2} \left(1 + \delta^2 + \eta^2 \sin^2 \theta \pm \text{Sign}(1 - \delta^2) \sqrt{(1 + \delta^2 + \eta^2 \sin^2 \theta)^2 - 4\delta^2} \right)$$

So basically,

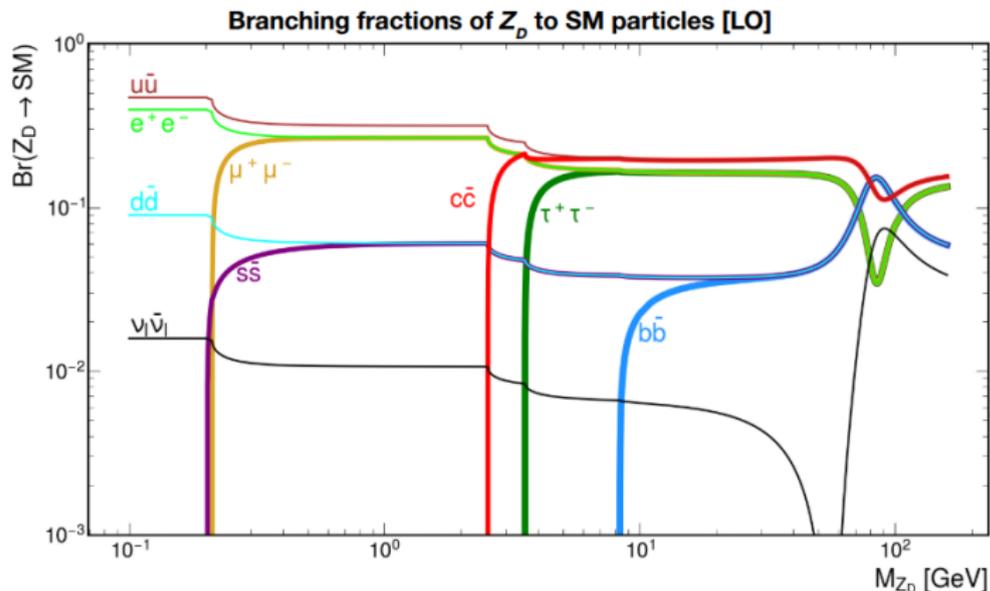
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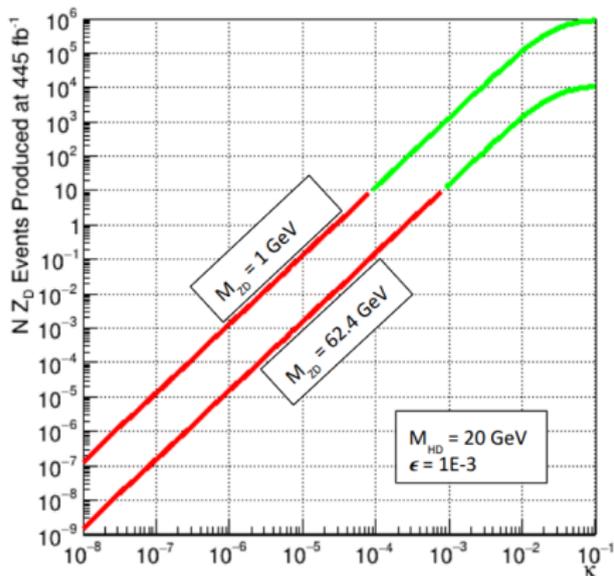
$$\Gamma(Z_D \rightarrow \bar{f}f) = \frac{N_c}{24\pi m_{Z_D}} \sqrt{1 - \frac{4m_f^2}{m_{Z_D}^2}} (m_{Z_D}^2 (g_L^2 + g_R^2) - m_f^2 (-6g_L g_R + g_L^2 + g_R^2))$$

HAHM basics

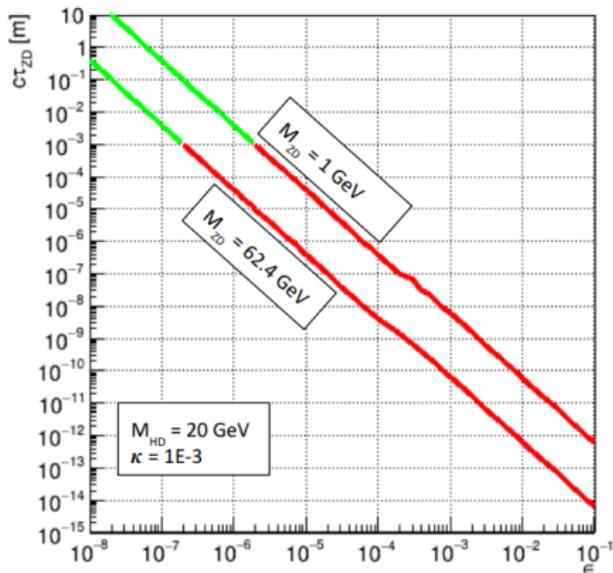


- Branching Ratios of Z_D to SM
 - How often does Z_D decay to each fermion pair based on M_{Z_D}
- $\text{Br}_{\text{families}}$ equal as $M_{Z_D} \geq 2m_f$
- Any lepton is good!

HAHM basics

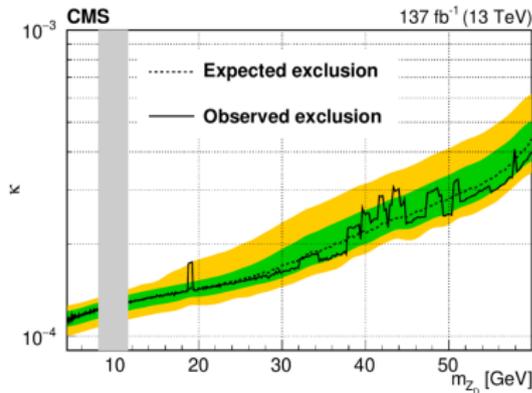
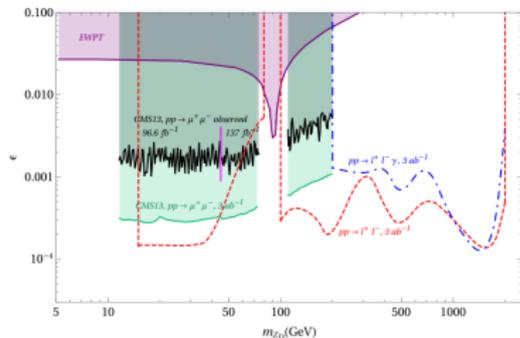


- $\text{Br}(h \rightarrow Z_D Z_D) \propto \kappa^2$
- $\sigma_{\text{Prod}}(h \rightarrow Z_D Z_D \rightarrow \text{llll}) \propto \kappa^2$
- We want κ large ($> 10^{-2}$)

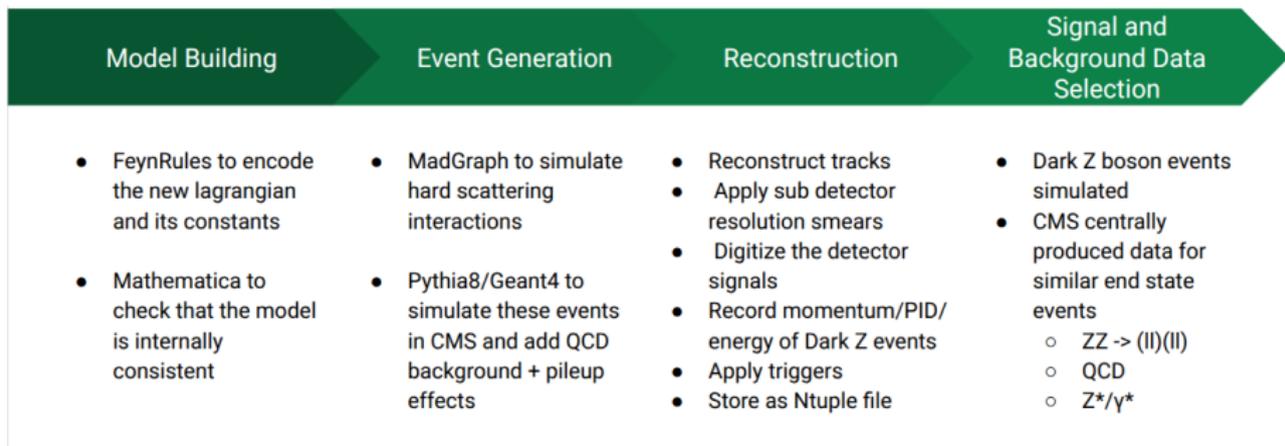


- $\Gamma(Z_D \rightarrow ff) \propto \epsilon^2$
- $\tau = \frac{\hbar}{\Gamma} \therefore c\tau \propto \epsilon^{-2}$
- We want ϵ small ($< 10^{-5}$)

- But *Why* HAHM?
- Improving LLP sensitivity greatly expands detectable ϵ
- Current Limits on ϵ and κ [10][11]
 - All Prompt and κ small to suppress $h \rightarrow h_D h_D$

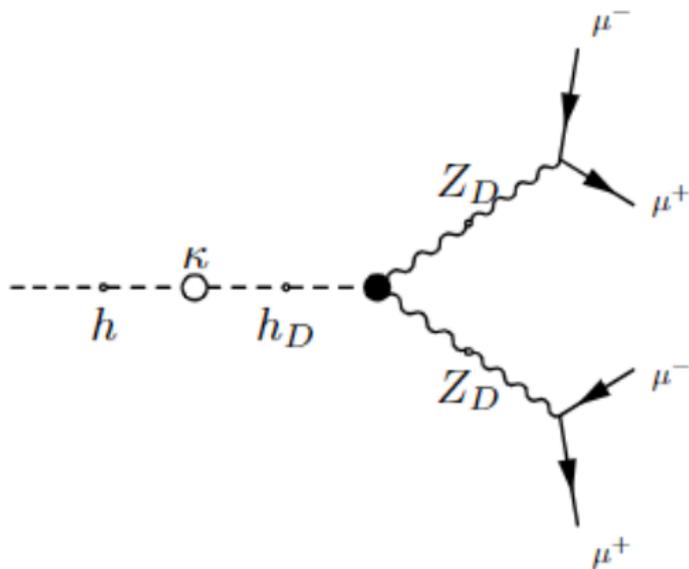


CMS simulation process



Four-Muon Final State

- Muons interact very weakly with detector material, so they can travel through the inner layers and reach the muon system.
- This makes muons very clean and easy to identify.
- The 125-GeV Higgs boson (h) is assumed to mix with a dark Higgs boson (h_D) through Higgs mixing κ
- Kinetic mixing ϵ could give Z_d small coupling to dimuons
- $h \rightarrow Z_D Z_D \rightarrow 2\mu^+ 2\mu^-$ is dominate if, $\epsilon \ll \kappa$, $m_{h_D} > \frac{m_h}{2}$, and $m_{Z_D} < \frac{m_h}{2}$



Four-Muon Final State – Signal Cutflow (50Gev, 50 mm Sample)

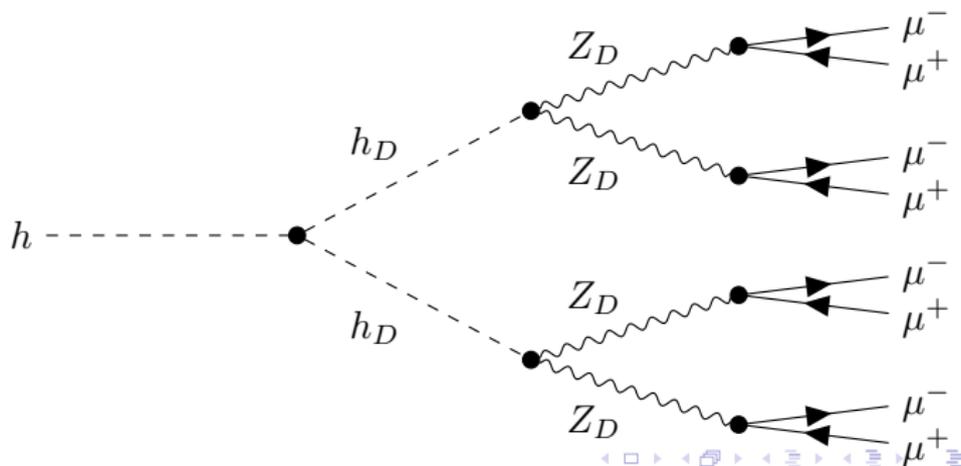
#	Cut	Events	Rel. Eff (%)	Total Eff (%)
0	All events	10000	–	100.0
1	GEN 1st μ $p_T > 24$, $ \eta < 2$	8639	86.4	86.4
2	GEN 2nd μ $p_T > 24$, $ \eta < 2$	6329	73.3	63.3
3	GEN 3rd μ $p_T > 10$, $ \eta < 2.4$	6114	96.6	61.1
4	GEN 4th μ $p_T > 10$, $ \eta < 2.4$	4513	73.8	45.1
5	GEN pixel/tracker fiducial	4042	89.6	40.4
6	RECO 1st μ $p_T > 24$, $ \eta < 2$ (rel. to Cut4)	4474	99.1	44.7
7	RECO 2nd μ $p_T > 24$, $ \eta < 2$	3936	88.0	39.4
8	RECO 3rd μ $p_T > 10$, $ \eta < 2.4$	3936	100.0	39.4
9	RECO 4th μ $p_T > 10$, $ \eta < 2.4$	3936	100.0	39.4
10	Good PV	3794	96.4	37.9
11	fitValid + prob < 0.01	2525	66.5	25.3
12	Valid pixel hit	2394	94.8	23.9
13	DY veto	2394	100.0	23.9
14	Iso_eachMuon < 2.3	1655	69.1	16.6
15	Signal HLT	1653	99.9	16.5
16	STA quality	1627	98.4	16.3
17	Mass consistency	1369	84.1	13.7

- Efficiency (vs cut5) $\approx 1369/4042 \sim 33\%$
- Efficiency (vs cut4) $\approx 1369/4513 \sim 30\%$

Eight Muon Endstate

- h can also decay into h_D
- h_D can decay into SM particles
- If $m_{h_D} > 2m_{Z_D}$, then $\text{Br}(h_D \rightarrow Z_D Z_D) \sim 100\%$
- Also, $\text{Br}(h \rightarrow h_D h_D) \sim \kappa^2$

- Pros
 - Nearly no background
 - h_D decay promptly to Z_D
- Cons
 - Limits range for M_{h_D} and M_{Z_D}
 - Easy to lose muons



Eight Muon Endstate – Pros

- Nearly background free process
 - Results normalized to Run2
 - $M_{Z_D} = 30\text{GeV}$, $M_{h_D} = 62\text{GeV}$
- h_D decays promptly to Z_D
 - h_D is short lived
 - Displacement only from Z_D

Background ($ZZ \rightarrow 4\ell$) – 181,525 Events

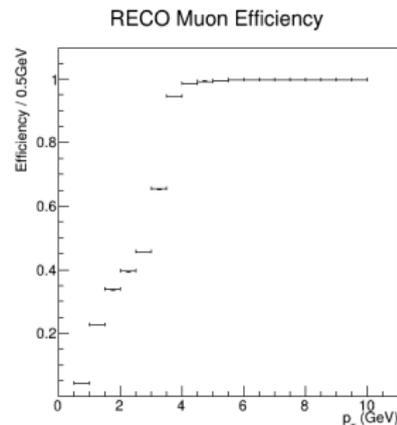
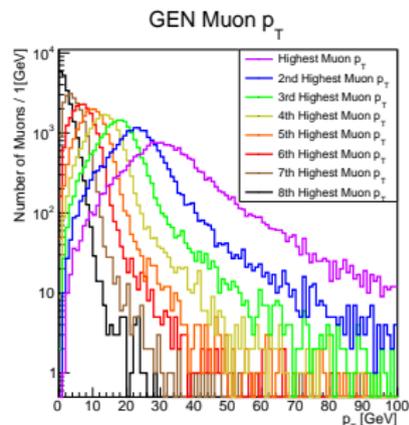
Num. RECO Muon	No Cuts	$p_T > 2\text{GeV}$	$p_T > 3\text{GeV}$	$p_T > 4\text{GeV}$	$p_T > 5\text{GeV}$
7	2735	27	16	6	3
8	1235	5	3	1	0
9 or more	909	34	19	7	4
7 or more	4879	66	38	14	7
7 or more	100%	1.34%	0.78%	0.29%	0.14%

Signal ($h \rightarrow h_D h_D$) – 20,000 Events

Num. RECO Muon	No Cuts	$p_T > 2\text{GeV}$	$p_T > 3\text{GeV}$	$p_T > 4\text{GeV}$	$p_T > 5\text{GeV}$
7	3727	3595	3479	3219	2823
8	8728	8485	8148	7454	6494
9 or more	1574	1562	1529	1453	1329
7 or more	14029	13642	13156	12126	10646
7 or more	100%	97.24%	93.78%	86.44%	75.89%

Eight Muon Endstate – Cons

- Limits mass of h_D and Z_D
 - Require $M_{h_D} < \frac{1}{2} M_h$
 - Require $M_{Z_D} < \frac{1}{2} M_{h_D}$
 - If either is not satisfied, this process is heavily suppressed
- Easy to lose muons
 - Energy split across 8 muons
 - Low energy muons cannot be reconstructed reliably
 - Long lived Z_D can decay outside of CMS
 - Want 8 muons in an event
 - Might only see 7 or less



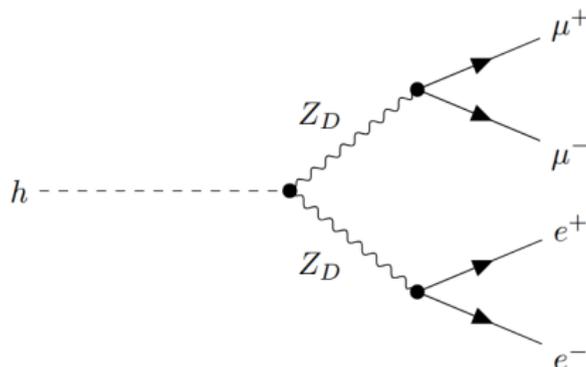
Eight Muon Endstate – Current and Future Work

- Generate Samples for Run 3
 - Samples spanning over:
 - $1\text{GeV} < M_{Z_D} < 30\text{GeV}$
 - $2M_{Z_D} < M_{h_D} < 62\text{GeV}$
 - $0\text{mm} < c\tau < 1000\text{mm}$
 - Samples are currently undergoing simulation
- Perform muon selection and muon pairing on samples
 - Make sure all muons selected are of good quality
 - Can also use this to reject some background
 - Currently working on improving this
- Ensure that background is nearly zero
 - Shown that $ZZ \rightarrow 4\ell$ is nearly completely removed
 - Need to show QCD background has the same behavior
 - Will then have to do data driven background estimation
- Discover new physics / place limits $\text{Br}(h \rightarrow h_D h_D)$

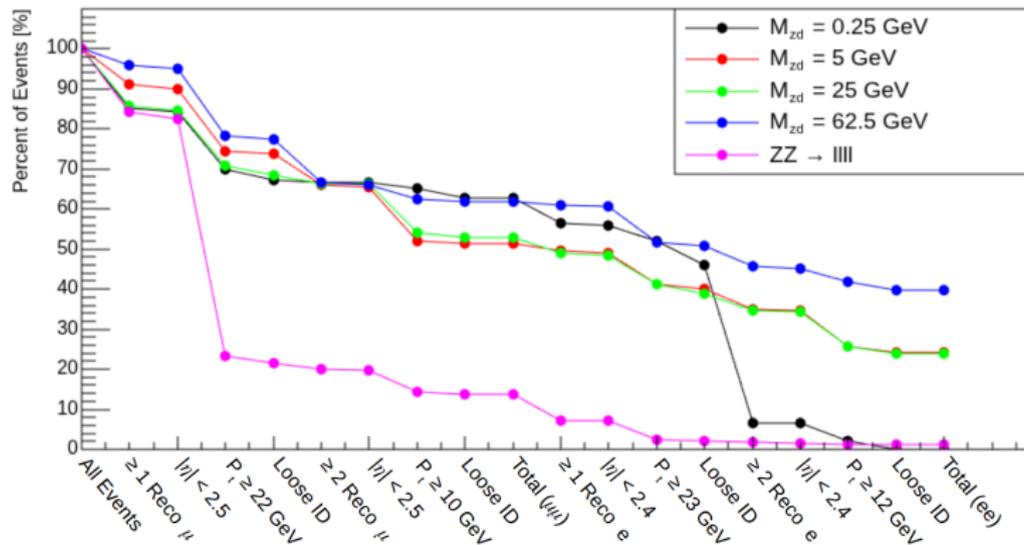
Two Electrons, Two Muons

- Include electrons in the analysis
- Set $M_{hd} = 130$ GeV
 - Not produced
- $M_{zd} \in [0.25, 62.5]$ GeV
- $c\tau \in [0, 1000]$ mm

- Pros
 - Increases signal events
 - Can still trigger on $(\mu\mu)$ events
- Cons
 - Electrons have large QCD background
 - if $M_{zd} < 10$ GeV, $(\Delta R_{ee} < 0.1)$!

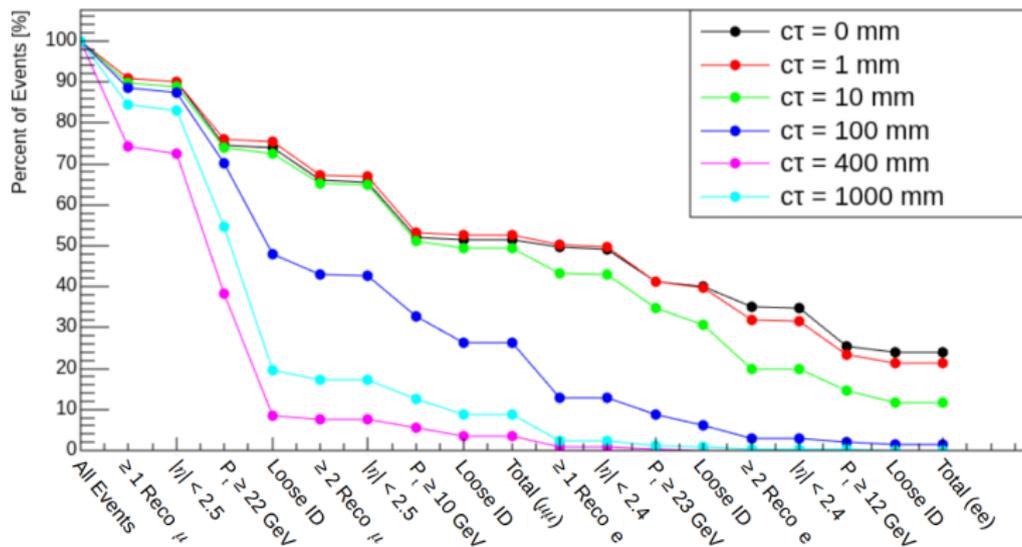


Two Electrons, Two Muons



- Apply selection criteria for "good" signal
 - Within certain η or P_t so its detectable (geometric acceptance)
- Background reduced to 1.2%

Two Electrons, Two Muons



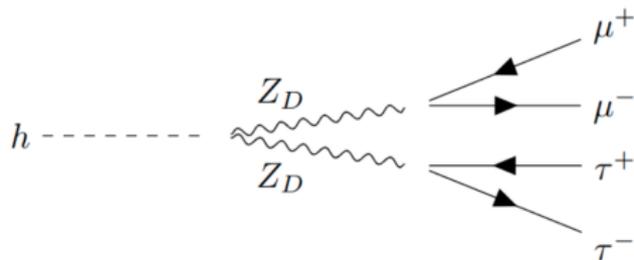
- We can also vary ct
 - Largest losses are due to identifying these leptons
- Large ct suffers with current CMS software

Two Electrons, Two Muons

Whats left to do?

- Optimize cuts applied to signal and background
- Apply offline triggers
- Better electron track identification
 - There are different kinds of lepton objects depending on the identification technique
- Currently reworking reconstructed lepton pairing
 - Naive: take the top two leptons with highest P_t
 - Now: Find (ee) $(\mu\mu)$ such that ΔM_{Inv} is smallest
- Potential to apply ML/NN to differentiate signal from background topology if kinematic separation isn't good enough
 - Especially for low mass dielectron resonance

Two Muons, Two Taus



τ leptons decay promptly and are never observed directly. They are only ever seen as their decay modes:

τ_e, τ_μ, τ_h

We choose the decay mode pairs with the highest branching fractions:

$\tau_h\tau_h, \tau_\mu\tau_h, \tau_e\tau_h, \tau_e\tau_\mu$

τ_e and τ_μ are reconstructed using standard electron and muon algorithms. τ_h requires special techniques...

Decay mode	Resonance	\mathcal{B} (%)
Leptonic decays		35.2
$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$		17.8
$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$		17.4
Hadronic decays		64.8
$\tau^- \rightarrow h^- \nu_\tau$		11.5
$\tau^- \rightarrow h^- \pi^0 \nu_\tau$	$\rho(770)$	25.9
$\tau^- \rightarrow h^- \pi^0 \pi^0 \nu_\tau$	$a_1(1260)$	9.5
$\tau^- \rightarrow h^- h^+ h^- \nu_\tau$	$a_1(1260)$	9.8
$\tau^- \rightarrow h^- h^+ h^- \pi^0 \nu_\tau$		4.8
Other		3.3

$$\text{Br}(\tau_h\tau_h) = 0.648^2 = 0.420$$

$$\text{Br}(\tau_\mu\tau_h) = 2 \cdot 0.174 \cdot 0.648 = 0.226$$

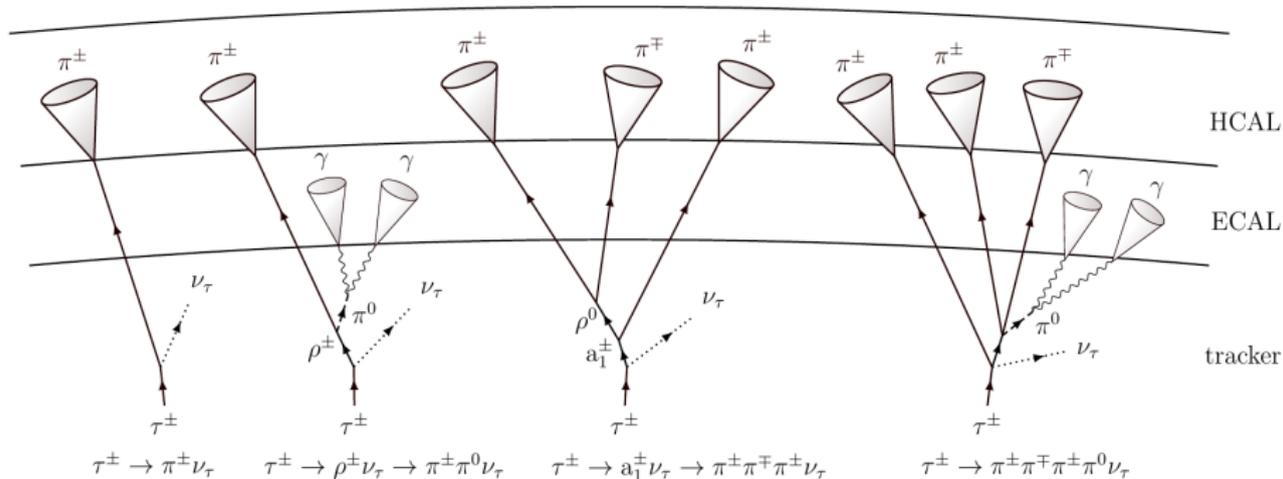
$$\text{Br}(\tau_e\tau_h) = 2 \cdot 0.178 \cdot 0.648 = 0.231$$

$$\text{Br}(\tau_e\tau_\mu) = 2 \cdot 0.178 \cdot 0.174 = 0.062$$

$$\text{Br}(\tau_e\tau_e) = 0.178^2 = 0.032$$

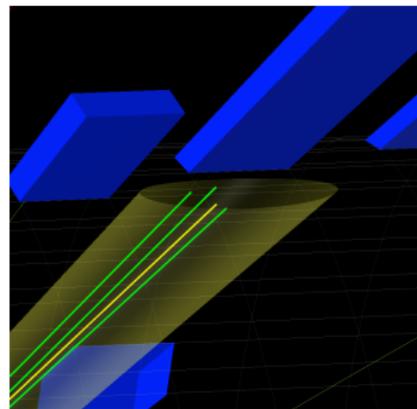
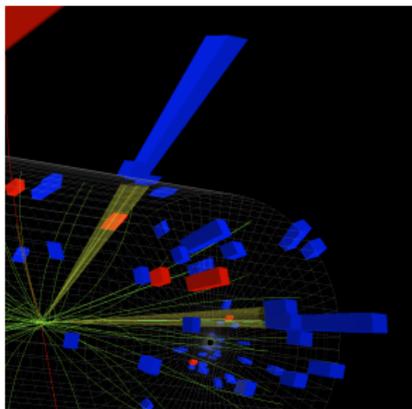
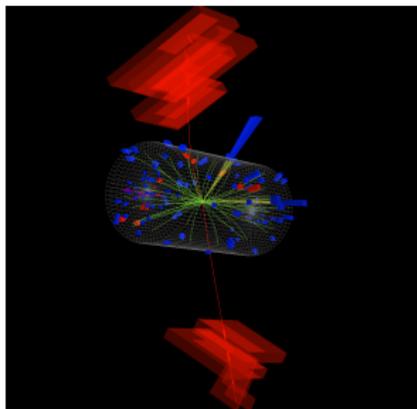
$$\text{Br}(\tau_\mu\tau_\mu) = 0.174^2 = 0.030$$

Hadronic Tau Reconstruction



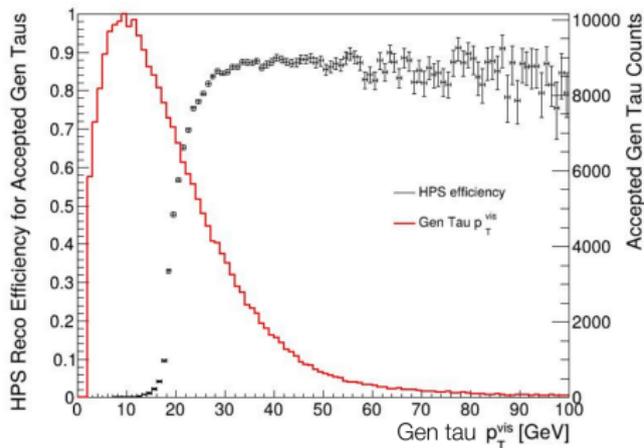
τ_h encompasses many different decay modes involving tracks, ECAL hits, and HCAL hits. The Hadrons Plus Strips (HPS) algorithm is the default.

A Good HPS Example



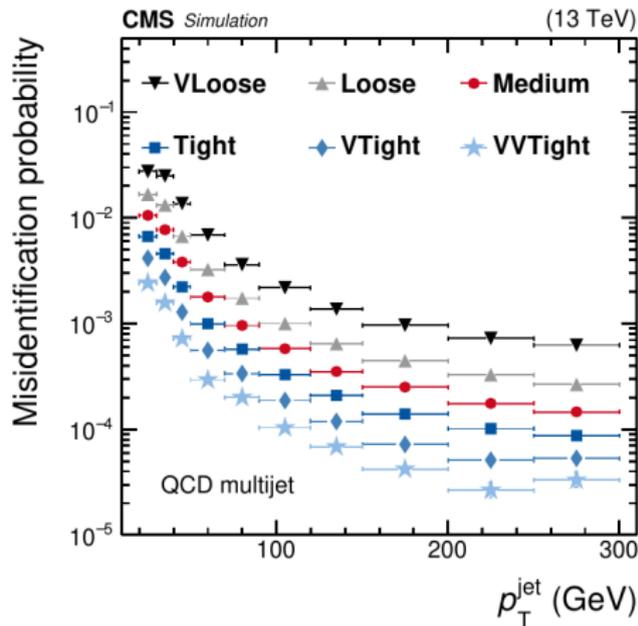
HPS has identified the three charged pion tracks (green) and the HCAL deposit (blue) as a τ_h reconstruction candidate (yellow).

Limitations of the HPS Algorithm



Our signal has relatively low momentum taus ($p_T < 20\text{GeV}$)

HPS has low efficiency and high fake rate in this region.



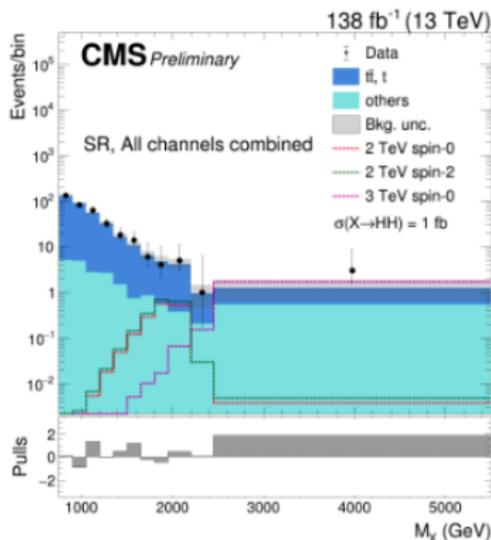
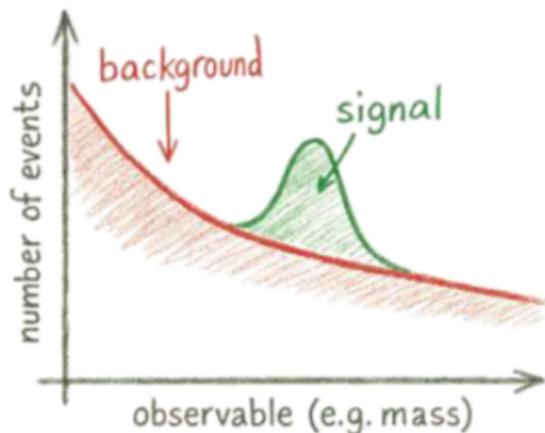
HPS performance is our biggest limitation.

We are doing a direct follow-up to a 2017 CMS study that uses HPS with a $\mu\mu\tau\tau$ final state. We are investigating alternative methods that have been developed since then:

- DeepTau NN applied on top of HPS candidates. Reduces fake rate but doesn't help efficiency.
- Scouting HPS developed for higher efficiency at low p_T , but for “scouting” datasets. Perhaps we can adapt it to our datasets.
- Custom NN to classify jet seeds.

Conclusions and Future plans

- Improving LLP sensitivity in CMS
 - We all must present our work to other BSM/LLP working groups
- Convince the community of our findings
 - $N_{Sig} > \sqrt{N_{bkgd}}$
- Release our analysis programs onto real CMS data (unblinding)
- Look for anomalies! [12]



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